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Journal of Nuclear Materials 283–287 (2000) 1414–1418

**journal of  
nuclear  
materials**

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# Improvement of the model for surface process of tritium release from lithium oxide

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## Abstract

Among the various tritium transport processes in lithium ceramics, the importance and the detailed mechanism of surface reactions remain to be elucidated. The dynamic adsorption and desorption model for tritium desorption from lithium ceramics, especially  $\text{Li}_2\text{O}$  was constructed. From the experimental results, it was considered that both  $\text{H}_2$  and  $\text{H}_2\text{O}$  are dissociatively adsorbed on  $\text{Li}_2\text{O}$  and generate  $\text{OH}^-$  on the surface. In the first model developed in 1994, it was assumed that either the dissociative adsorption of  $\text{H}_2$  or  $\text{H}_2\text{O}$  on  $\text{Li}_2\text{O}$  generates two  $\text{OH}^-$  on the surface. However, recent calculation results show that the generation of one  $\text{OH}^-$  and one  $\text{H}^-$  is more stable than that of two  $\text{OH}^-$ s by the dissociative adsorption of  $\text{H}_2$ . Therefore, assumption of  $\text{H}_2$  adsorption and desorption in the first model is improved and the tritium release behavior from  $\text{Li}_2\text{O}$  surface is evaluated again by using the improved model. The tritium residence time on the  $\text{Li}_2\text{O}$  surface is calculated using the improved model, and the results are compared with the experimental results. The calculation results using the improved model agree well with the experimental results than those using the first model. © 2000 Elsevier Science B.V. All rights reserved.

## 1. Introduction

Estimation of the precise tritium inventory in the fusion reactor blanket and the chemical form of the released tritium are essential for the effective recovery of tritium in the fusion reactor fuel cycle. To estimate the tritium inventory in the blanket system and the chemical form of the released tritium, a thorough understanding of tritium release behavior from breeding materials is required. In sintered solid breeding materials, tritium release is composed of a series of migration steps. Among them, surface reactions are known to be important.

Computer modeling is an important tool for clarifying the mechanism of surface reaction in the tritium release process and for estimating tritium inventory in the blanket and the chemical form of the released tritium [1–5]. The DAD (the dynamic adsorption and desorp-

tion) model was developed in 1994 to describe the surface processes of tritium release from  $\text{Li}_2\text{O}$  and was used to estimate the tritium inventory on the surface and the chemical form of the released tritium [6]. From the results of in situ tritium release experiments (TTTex), it was clarified that tritium release rate from the  $\text{Li}_2\text{O}$  surface is affected both by the swamping effect with  $\text{H}_2$  and  $\text{H}_2\text{O}$  and by the surface oxygen potential effect [7]. To describe both effects, the first DAD model developed in 1994 assumed that the surface is covered with hydroxyl ions, oxygen ions and vacancies on interstitial and lattice sites, that is, the swamping effect is considered to be described as the fraction of hydroxyl ions and the oxygen potential effect as the unbalance between the fraction of oxygen ions and vacancies.

In the first DAD model, the fractions of these species are determined from the partial pressures of  $\text{H}_2$  and  $\text{H}_2\text{O}$  in the sweep gas by adsorption and desorption reactions of  $\text{H}_2$  and  $\text{H}_2\text{O}$ . Tritium is released by three processes: as HTO by recombination of  $\text{OH}^-$  and  $\text{OT}^-$  on the surface, as HT by recombination of  $\text{OH}^-$  and  $\text{OT}^-$  at low oxygen potential or when these surface groups are surrounded by ‘active’ oxygen vacancies, and as HT by exchange reaction of surface  $\text{OT}^-$  with  $\text{H}_2$  in

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the gaseous phase. However, by the recent calculation results [8], it is clarified that the assumption of the  $H_2$  adsorption and desorption processes in the DAD model is not suitable for describing the tritium release surface processes.

In the present study, the DAD model is improved by introducing the new assumption for  $H_2$  adsorption and desorption process founded on the recent calculation results. The dependencies of the fraction of surface adsorbed species and the surface tritium residence time on the sweep gas chemical form are calculated, using the ‘improved’ DAD model, and the calculation results are compared with the experimental results.

## 2. Improvement of the model

In the DAD model, the (110) surface of  $Li_2O$  was selected for modeling. On the ideal (110) surface, as shown in the top three rows of Fig. 1,  $O^{2-}$  on the  $O^{2-}$  lattice sites and empty site ( $E_O$ ) located on the interstitial sites of the  $O^{2-}$  sub-lattice are arranged regularly among the  $Li^+$  lattice. These two kinds of sites will become the sites for  $OH^-$  adsorption ( $O_{ads}$  site in Fig. 1). Strictly speaking,  $E_O$  (empty site on interstitial sites of  $O^{2-}$  sub-lattice) and  $V_O$  (vacancy on  $O^{2-}$  lattice sites) are energetically different. However, it is considered that the surface species,  $OH^-$  and  $O^{2-}$ , can occupy both the lattice and the interstitial sites because  $OH^-$  and  $O^{2-}$  can move easily on the surface. Therefore, at the surface, it is considered that there is no significant difference in the energetics between the  $E_O$  on the interstitial site and the  $V_O$  on the lattice site. Therefore, in the DAD model, for simplification of the model calculation,  $E_O$  and  $V_O$  are treated equally as a kind of oxygen vacancy at the surface. The symbol  $V_S$  is used for such an oxygen vacancy at the surface, including both  $E_O$  and  $V_O$ .

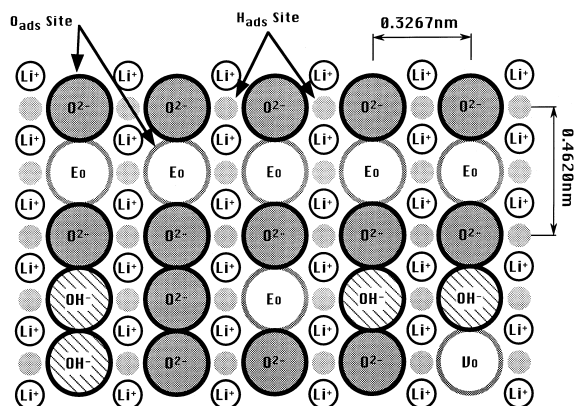


Fig. 1. The (110) surface of the  $Li_2O$  crystal.  $OH^-$ s are absorbed partially.

In the first DAD model,  $H_2O$  and  $H_2$  adsorption and desorption processes were assumed to be as follows [6]:  $H_2O$  is dissociatively adsorbed on two adjacent  $O_{ads}$  sites occupied by one  $V_S$  and one  $O^{2-}$ , and then produces two  $OH^-$ s.  $H_2$  is dissociatively adsorbed on two adjacent  $O_{ads}$  sites occupied by two  $O^{2-}$ s and produces two  $OH^-$ s.  $H_2O$  is desorbed from two adjacent  $O_{ads}$  sites occupied by two  $OH^-$ s by recombination process leaving  $V_S$  and  $O^{2-}$ . When the surface oxygen potential is very low,  $V_S$  existing on the adjacent sites of  $OH^-$ s are considered to affect the nature of  $OH^-$ s. Such  $V_S$  which we named active oxygen vacancy ( $V_O^a$ ) was considered to attract the oxygen atoms of  $OH^-$ s to the surface and the recombined  $OH^-$ s were considered to be desorbed as  $H_2$ .

The calculation results of the surface tritium residence time by using the first DAD model showed good agreement with the results of the in situ experiments [6]. However, some issues were pointed out in the assumptions in the first model. It seemed that electric balance was not maintained before and after the  $H_2$  adsorption,  $H_2 + 2O^{2-} \rightarrow 2OH^-$ . In addition, the definition of active oxygen vacancy which plays an important role in  $H_2$  desorption was not clear. These issues remained to be improved founded on the detailed study on the adsorption and desorption mechanism of  $H_2$ .

Recently, Taniguchi et al. performed the quantum chemical calculations to simulate the dissociative adsorption of  $H_2$  on the  $Li_2O$  surface [8]. They estimated the heat of  $H_2$  dissociative adsorption on the three types of possible adsorption sites. The three sites were as follows: (1) two dissociative H atoms are adsorbed on two  $O^{2-}$  and produce two hydroxyl ions, same as the assumption in the first DAD model, (2) one dissociative H atom is adsorbed on  $O^{2-}$  to produce  $OH^-$  and another dissociative H atom is adsorbed on the nearest  $Li^+$ , and (3) one dissociative H atom is adsorbed on  $O^{2-}$  to produce  $OH^-$  and another dissociative H atom is adsorbed on the site existing between the two nearest lithium ions ( $H_{ads}$  site in Fig. 1). Their calculation results showed that the last adsorption site is the most stable site for dissociative adsorption of  $H_2$  on  $Li_2O$ .

Based on these calculation results, the alternation of the assumption of  $H_2$  adsorption and desorption mechanism in the first DAD model is performed, and the improved DAD model was constructed. In the first model, only  $O_{ads}$  sites in Fig. 1 were considered as the dissociative adsorption sites for  $H_2$  and  $H_2O$ . In the improved model,  $H_{ads}$  sites in Fig. 1 are also considered as the dissociative adsorption sites for  $H_2$ .

The assumption of the  $H_2$  and  $H_2O$  adsorption and desorption mechanisms in the improved model is as follows:  $H_2O$  adsorption and desorption are assumed to occur on the  $O_{ads}$  sites in the same process as assumed in the first model.  $H_2$  is assumed to be dissociatively adsorbed on  $O^{2-}$  and adjacent  $E_H$  ( $H_{ads}$  site which is not

occupied by  $H^-$  and produces  $OH^-$  and  $H^-$ . When  $OH^-$  does not have any neighboring  $OH^-$  and has adjacent  $H^-$  on the nearest  $H_{\text{ads}}$  site,  $H_2$  is assumed to be desorbed by recombination process of  $OH^-$  and  $H^-$  and leaves  $O^{2-}$  and  $E_H$ .

### 3. Calculation

The  $H_2O$  adsorption rate is calculated as follows: (1) The probability of  $O_{\text{ads}}$  site occupied by a  $V_S$  which has a neighboring  $O^{2-}$  or by  $O^{2-}$  which has a neighboring  $V_S$  is  $\theta_V[1 - (1 - \theta_O)^2] + \theta_O[1 - (1 - \theta_V)^2]$ , where  $\theta_V$  and  $\theta_O$  are the surface coverages of  $V_S$  and  $O^{2-}$  on  $O_{\text{ads}}$  sites, respectively. (2) As the  $H_2O$  adsorption rate should be proportional to the  $H_2O$  pressure in the gaseous phase near the surface, the  $H_2O$  adsorption rate on the pair of sites occupied by a  $V_S$  and  $O^{2-}$  is described as  $K_{H_2O}^a \{\theta_V[1 - (1 - \theta_O)^2] + \theta_O[1 - (1 - \theta_V)^2]\} P_{H_2O}$ , where  $K_{H_2O}^a$  (molecules  $\text{Pa}^{-1} \text{s}^{-1} \text{site}^{-1}$ ) is the rate constant  $H_2O$  adsorption and  $P_{H_2O}$  (Pa) the partial pressure of  $H_2O$ . (3) The number of the  $O_{\text{ads}}$  sites on the  $Li_2O$  surface is equal to  $N_{SO}A$ , where  $N_{SO}$  (sites  $\text{m}^{-2}$ ) is the area density of  $O_{\text{ads}}$  sites and  $A$  ( $\text{m}^2$ ) is the surface area of  $Li_2O$ . Then, the overall  $H_2O$  adsorption rate  $R_{H_2O}^a$  (molecules  $\text{s}^{-1}$ ) is

$$R_{H_2O}^a = K_{H_2O}^a \{\theta_V[1 - (1 - \theta_O)^2] + \theta_O[1 - (1 - \theta_V)^2]\} P_{H_2O} N_{SO} A. \quad (1)$$

The overall  $H_2$  adsorption rate  $R_{H_2}^a$  (molecules  $\text{s}^{-1}$ ) is calculated in the same way as the  $H_2O$  adsorption,

$$R_{H_2}^a = K_{H_2}^a \theta_O [1 - (1 - \theta_{EH})^2] P_{H_2} N_{SO} A, \quad (2)$$

where  $K_{H_2}^a$  (molecules  $\text{Pa}^{-1} \text{s}^{-1} \text{site}^{-1}$ ) is the rate constant of  $H_2$  adsorption,  $\theta_{EH}$  the surface coverage of  $E_H$  on  $H_{\text{ads}}$  sites and  $P_{H_2}$  (Pa) is the partial pressure of  $H_2$ .

Since the probability of one pair of neighboring  $O_{\text{ads}}$  sites being occupied by two  $OH^-$  is  $\theta_{OH}[1 - (1 - \theta_{OH})^2]$ , the overall  $H_2O$  desorption rate  $R_{H_2O}^d$  (molecules  $\text{s}^{-1}$ ) is described as

$$R_{H_2O}^d = K_{H_2O}^d \theta_{OH} [1 - (1 - \theta_{OH})^2] N_{SO} A, \quad (3)$$

where  $K_{H_2O}^d$  (molecules  $\text{s}^{-1} \text{site}^{-1}$ ) is the rate constant of  $H_2O$  desorption and  $\theta_{OH}$  the surface fraction of  $OH^-$  on  $O_{\text{ads}}$  sites.

The overall  $H_2$  desorption rate  $R_{H_2}^d$  (molecules  $\text{s}^{-1}$ ) is calculated in the same way as the  $H_2O$  desorption,

$$R_{H_2}^d = K_{H_2}^d \theta_{OH} [1 - (1 - \theta_H)^2] (1 - \theta_{OH})^2 N_{SO} A, \quad (4)$$

where  $K_{H_2}^d$  (molecules  $\text{s}^{-1} \text{site}^{-1}$ ) is the rate constant of  $H_2$  desorption and  $\theta_H$  is the surface coverage of  $H^-$  on  $H_{\text{ads}}$  sites.

The relation among the surface coverages is expressed by

$$\begin{aligned} \theta_O + \theta_V + \theta_{OH} &= 1, & \theta_H + \theta_{EH} &= 1, & 0 < \theta_O < 1, \\ 0 < \theta_V < 1, & 0 < \theta_{OH} < 1, & & & (5) \\ 0 < \theta_H < 1, & 0 < \theta_{EH} < 1. & & & \end{aligned}$$

On the ideal (110) surface of  $Li_2O$  which absorbs neither hydrogen nor hydroxyl groups, the relations of  $\theta_V = \theta_O$  and  $\theta_H = 0$  are established. In the improved model assumption,  $H^-$  on  $H_{\text{ads}}$  sites increases and  $O^{2-}$  on  $O_{\text{ads}}$  sites decreases by the adsorption of  $H_2$  molecule. Also,  $H^-$  on  $H_{\text{ads}}$  site decreases and  $O^{2-}$  on  $O_{\text{ads}}$  site increases by the desorption of  $H_2$  molecule. The relation of  $\theta_V = \theta_O$  and the value of  $\theta_H$  does not change by the adsorption and desorption of  $H_2O$ . As shown in Fig. 1,  $N_{SH}$  is equal to  $N_{SO}$ . Therefore, the relation of  $\theta_H + \theta_O = \theta_V$  is established on the ideal (110) surface in the  $H_2$  and  $H_2O$  adsorption and desorption system.

When the equilibrium state of  $H_2O$  adsorption and desorption is established, from Eqs. (1) and (3),

$$K_{H_2O}^a \{\theta_V[1 - (1 - \theta_O)^2] + \theta_O[1 - (1 - \theta_V)^2]\} P_{H_2O} - K_{H_2O}^d \theta_{OH} [1 - (1 - \theta_{OH})^2] = 0 \quad (6)$$

When the equilibrium state of  $H_2$  adsorption and desorption is established, from Eqs. (2) and (4),

$$K_{H_2}^a \theta_O [1 - (1 - \theta_{EH})^2] P_{H_2} - K_{H_2}^d \theta_{OH} [1 - (1 - \theta_H)^2] \times (1 - \theta_{OH})^2 = 0 \quad (7)$$

Tritium generated in  $Li_2O$  is considered to exist as  $OT^-$  on  $O_{\text{ads}}$  sites or  $T^-$  on  $H_{\text{ads}}$  sites by exchange reaction. In the improved model, it is assumed that the ratio of  $OT^-/T^-$  on  $Li_2O$  surface is equal to that of  $OH^-/H^-$  on the surface.

When the surface density of tritium is much less than that of hydrogen, tritium on the surface is considered to be released as HTO by recombination of surface  $OH^-$  and  $OT^-$ , and HT by recombination of surface  $OH^-$  and  $T^-$ , or  $OT^-$  and  $H^-$ , which are the similar processes to the desorption of  $H_2O$  or  $H_2$ . Moreover, from the in situ experimental results, HT is also considered to be released by the exchange reaction between  $H_2$  in the sweep gas and  $OT^-$  or  $T^-$  on the surface [7]. In this improved model, it is assumed that the HT is released by the exchange reaction from  $OT^-$  which has no  $OH^-$  on its adjacent  $O_{\text{ads}}$  sites and no  $H^-$  on its adjacent  $H_{\text{ads}}$  sites, and from  $T^-$  which has no  $OH^-$  on its adjacent  $O_{\text{ads}}$  sites. Therefore, the total tritium release rate  $R_T$  is described as

$$\begin{aligned} R_T &= K_{HTO}^r \theta_{OT} [1 - (1 - \theta_{OH})^2] N_{SO} A \\ &+ K_{HT}^r \{\theta_{OT} [1 - (1 - \theta_H)^2] (1 - \theta_{OH})^2 \\ &+ \theta_T [1 - (1 - \theta_{OH})^2]\} N_{SO} A \\ &+ K_{HT}^e [\theta_{OT} (1 - \theta_{OH}^2) N_{SO} A \\ &+ \theta_T (1 - \theta_{OH}^2) (1 - \theta_H^2) N_{SH} A] P_{H_2}, \end{aligned} \quad (8)$$

where  $K_{\text{HTO}}^r$  (molecules  $\text{s}^{-1}$  site $^{-1}$ ),  $K_{\text{HT}}^r$  (molecules  $\text{s}^{-1}$  site $^{-1}$ ) and  $K_{\text{HT}}^s$  (molecules  $\text{Pa}^{-1}$   $\text{s}^{-1}$  site $^{-1}$ ) are the release rate constants of HTO by recombination, HT by recombination and HT by exchange reaction, respectively,  $R_T$  (molecules  $\text{s}^{-1}$ ) the overall tritium release rate,  $\theta_{\text{OT}^-}$  is the surface fraction of  $\text{OT}^-$  on  $\text{O}_{\text{ads}}$  sites, and  $\theta_T$  is the surface fraction of  $\text{T}^-$  on  $\text{H}_{\text{ads}}$  sites. The surface tritium residence time  $\tau_T$  (s) is

$$\tau_T = (\theta_{\text{OT}^-} N_{\text{SO}} + \theta_T N_{\text{SH}}) A / R_T. \quad (9)$$

The kinetic constants in the improved model are determined by comparing with in situ and out-of-pile experiments at 773 K.

When no  $\text{H}_2$  exists in the gaseous phase, the relation of  $\theta_{\text{OH}} = \theta_{\text{O}}$  and  $\theta_{\text{H}} = 0$  is established, and Eqs. (6) and (8) are simplified.  $K_{\text{HTO}}^r$  is considered to be nearly equal to  $K_{\text{H}_2\text{O}}^d$ . From Eq. (9) and the simplified Eqs. (6) and (8),  $K_{\text{H}_2\text{O}}^d$  ( $K_{\text{HTO}}^r$ ) and  $K_{\text{H}_2\text{O}}^a$  were calculated to be  $1.7 \times 10^{-2}$  (molecules  $\text{s}^{-1}$  site $^{-1}$ ) and  $1.9 \times 10^{-5}$  (molecules  $\text{Pa}^{-1}$   $\text{s}^{-1}$  site $^{-1}$ ), respectively, by using experimental data for the tritium surface residence time for the  $\text{He} + \text{H}_2\text{O}$  sweep gas [7] and the amount of  $\text{H}_2\text{O}$  adsorption on  $\text{Li}_2\text{O}$  [9].

It is not easy to determine  $K_{\text{H}_2}^d$  and  $K_{\text{H}_2}^a$  only from in situ or out-of-pile experiments because at actual conditions the experimental result is inevitably affected by co-existing  $\text{H}_2\text{O}$ . In the improved model, their values are determined from the balance between  $\theta_{\text{O}}$  and  $\theta_{\text{V}}$ , which is the similar process of evaluation in the first DAD model. For the condition in which  $P_{\text{H}_2\text{O}}$  is much larger than  $P_{\text{H}_2}$ ,  $\theta_{\text{V}}$  is considered to be nearly equal to  $\theta_{\text{O}}$ , because the surface balance is mainly determined by the adsorption and desorption of  $\text{H}_2\text{O}$ . If the effect of  $\text{H}_2$  becomes significant, imbalance occurs between  $\theta_{\text{V}}$  and  $\theta_{\text{O}}$ . In the improved model, it has been assumed that the imbalance between  $\theta_{\text{V}}$  and  $\theta_{\text{O}}$  begins at  $P_{\text{H}_2\text{O}} = 100$  Pa and  $P_{\text{H}_2} = 1$  Pa, which is the same assumption as in the first model [6]. It led to  $K_{\text{H}_2}^d = 2.0 \times 10^{-4}$  (molecules  $\text{s}^{-1}$  site $^{-1}$ ) and  $K_{\text{H}_2}^a = 9.5 \times 10^{-7}$  (molecules  $\text{Pa}^{-1}$   $\text{s}^{-1}$  site $^{-1}$ ) by using experimental data for the tritium surface residence time for the  $\text{He} + \text{H}_2 + \text{H}_2\text{O}$  sweep gas [10].

$K_{\text{HT}}^r$  is considered to be nearly equal to  $K_{\text{H}_2}^d$ , so  $K_{\text{HT}}^r = 2.0 \times 10^{-4}$  (molecules  $\text{s}^{-1}$  site $^{-1}$ ).

$K_{\text{HT}}^s$  is calculated from the results of in situ experiments of the sample at low oxygen potential, where tritium surface residence time was found to be inversely proportional to the hydrogen pressure, to be  $6.0 \times 10^{-5}$  (molecules  $\text{s}^{-1}$  site $^{-1}$ ).

#### 4. Results and discussions

Fig. 2 shows the surface coverages of  $\text{OH}^-$ ,  $\text{O}^{2-}$  and  $\text{V}_\text{S}$  on  $\text{O}_{\text{ads}}$  sites. In the calculation results using the first DAD model, large imbalance appears between  $\theta_{\text{O}}$  and

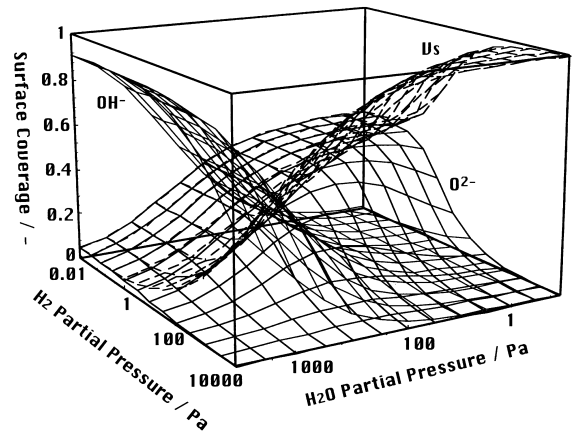


Fig. 2. Calculated surface coverages on  $\text{O}_{\text{ads}}$  sites.

$\theta_{\text{V}}$  even in the region for which  $P_{\text{H}_2\text{O}}$  is much larger than  $P_{\text{H}_2}$ . However, using the improved DAD model, no imbalance appears between  $\theta_{\text{O}}$  and  $\theta_{\text{V}}$  in the region for which  $P_{\text{H}_2\text{O}}$  is much larger than  $P_{\text{H}_2}$  (see Fig. 2). This result shows that the assumed mechanism of  $\text{H}_2$  adsorption and desorption in the improved model is more adequate for evaluating the tritium behavior on  $\text{Li}_2\text{O}$  surface than that in the first model.

Fig. 3 shows the calculated surface tritium residence time using the improved model compared with the experimental results from TTTEX [10]. Characteristics of model calculation are as follows: (1)  $\tau_T$  is inversely proportional to  $P_{\text{H}_2}$  at low  $P_{\text{H}_2\text{O}}$ , (2) inversely proportional to the square root of  $P_{\text{H}_2\text{O}}$  at low  $P_{\text{H}_2}$ , and (3) has a ridge on the line where  $P_{\text{H}_2\text{O}}$  is nearly equal to  $P_{\text{H}_2}$ . These characteristics of the calculated results agree well with the experimental results, and the estimated values of  $\tau_T$

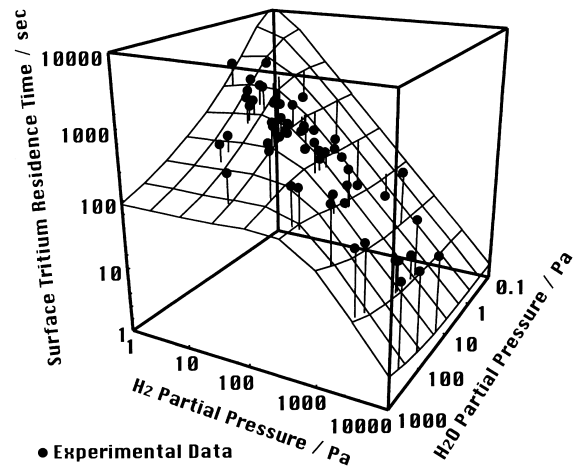


Fig. 3. Calculated surface tritium residence time compared with the experimental data from TTTEX [10].

agree more satisfactorily with the experimental results than the first DAD model, especially in the regions of  $P_{\text{H}_2} \leq 500$  Pa and  $P_{\text{H}_2\text{O}} \leq 100$  Pa [6]. However, there is a tendency for the difference between the experimental and calculated to become larger in a higher  $P_{\text{H}_2}$  region, same as with the first DAD model. A possible explanation for this tendency is the large experimental error due to the small residence time, or that the equilibrium on the surface is not established in higher  $P_{\text{H}_2}$  region.

## 5. Conclusion

The DAD model, which explains both the effects of swamping with hydrogen isotopes and the surface oxygen potential on tritium release from the  $\text{Li}_2\text{O}$  surface, is improved according to the recent calculation results performed by Taniguchi et al. The results of this improved DAD model show better agreement with the data from in situ experiments (TTTE<sub>x</sub>) than the first DAD model.

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